

LOW ENERGY SOLUTIONS FOR THE SEMICLASSICAL LIMIT OF SCHROEDINGER MAXWELL SYSTEMS

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ABSTRACT. We show that the number of solutions of Schroedinger Maxwell system on a smooth bounded domain $\Omega \subset \mathbb{R}^3$, depends on the topological properties of the domain. In particular we consider the Lusternik-Schnirelmann category and the Poincaré polynomial of the domain.

Dedicated to our friend Bernhard

1. INTRODUCTION

Given real numbers $q > 0, \omega > 0$ we consider the following Schroedinger Maxwell system on a smooth bounded domain $\Omega \subset \mathbb{R}^3$.

$$(1) \quad \begin{cases} -\varepsilon^2 \Delta u + u + \omega uv = |u|^{p-2}u & \text{in } \Omega \\ -\Delta v = qu^2 & \text{in } \Omega \\ u, v = 0 & \text{on } \partial\Omega \end{cases}$$

This paper deals with the semiclassical limit of the system (1), i.e. it is concerned with the problem of finding solutions of (1) when the parameter ε is sufficiently small. This problem has some relevance for the understanding of a wide class of quantum phenomena. We are interested in the relation between the number of solutions of (1) and the topology of the bounded set Ω . In particular we consider the Lusternik Schnirelmann category $\text{cat } \Omega$ of Ω in itself and its Poincaré polynomial $P_t(\Omega)$.

Our main results are the following.

Theorem 1. *Let $4 < p < 6$. For ε small enough there exist at least $\text{cat}(\Omega)$ positive solutions of (1).*

Theorem 2. *Let $4 < p < 6$. Assume that for ε small enough all the solutions of problem (1) are non-degenerate. Then there are at least $2P_1(\Omega) - 1$ positive solutions.*

Schroedinger Maxwell systems recently received considerable attention from the mathematical community. In the pioneering paper [9] Benci and Fortunato studied system (1) when $\varepsilon = 1$ and without nonlinearity. Regarding the system in a semiclassical regime Ruiz [18] and D'Aprile-Wei [11] showed the existence of a family of radially symmetric solutions respectively for $\Omega = \mathbb{R}^3$ or a ball. D'Aprile-Wei [12] also proved the existence of clustered solutions in the case of a bounded domain Ω in \mathbb{R}^3 .

Date: March 28, 2013.

2010 Mathematics Subject Classification. 35J60, 35J20, 58E30, 81V10.

Key words and phrases. Schroedinger-Maxwell systems, Lusternik Schnirelman category, Morse Theory.

Recently, Siciliano [19] relates the number of solution with the topology of the set Ω when $\varepsilon = 1$, and the nonlinearity is a pure power with exponent p close to the critical exponent 6. Moreover, in the case $\varepsilon = 1$, many authors proved results of existence and non existence of solution of (1) in presence of a pure power nonlinearity $|u|^{p-2}u$, $2 < p < 6$ or more general nonlinearities [1, 2, 3, 4, 10, 14, 15, 17, 20].

In a forthcoming paper [13], we aim to use our approach to give an estimate on the number of low energy solutions for Klein Gordon Maxwell systems on a Riemannian manifold in terms of the topology of the manifold and some information on the profile of the low energy solutions.

In the following we always assume $4 < p < 6$.

2. NOTATIONS AND DEFINITIONS

In the following we use the following notations.

- $B(x, r)$ is the ball in \mathbb{R}^3 centered in x with radius r .
- The function $U(x)$ is the unique positive spherically symmetric function in \mathbb{R}^3 such that

$$-\Delta U + U = U^{p-1} \text{ in } \mathbb{R}^3$$

we remark that U and its first derivative decay exponentially at infinity.

- Given $\varepsilon > 0$ we define $U_\varepsilon(x) = U\left(\frac{x}{\varepsilon}\right)$.
- We denote by $\text{supp } \varphi$ the support of the function φ .
- We define

$$m_\infty = \inf_{\int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx = |v|_{L^p(\mathbb{R}^3)}^p} \frac{1}{2} \int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx - \frac{1}{p} |v|_{L^p(\mathbb{R}^3)}^p$$

- We also use the following notation for the different norms for $u \in H_g^1(M)$:

$$\begin{aligned} \|u\|_\varepsilon^2 &= \frac{1}{\varepsilon^3} \int_M \varepsilon^2 |\nabla u|^2 + u^2 dx & |u|_{\varepsilon,p}^p &= \frac{1}{\varepsilon^3} \int_\Omega |u|^p dx \\ \|u\|_{H_0^1}^2 &= \int_\Omega |\nabla u|^2 dx & |u|_p^p &= \int_\Omega |u|^p dx \end{aligned}$$

and we denote by H_ε the Hilbert space $H_0^1(\Omega)$ endowed with the $\|\cdot\|_\varepsilon$ norm.

Definition 3. Let X a topological space and consider a closed subset $A \subset X$. We say that A has category k relative to X ($\text{cat}_M A = k$) if A is covered by k closed sets A_j , $j = 1, \dots, k$, which are contractible in X , and k is the minimum integer with this property. We simply denote $\text{cat } X = \text{cat}_X X$.

Remark 4. Let X_1 and X_2 be topological spaces. If $g_1 : X_1 \rightarrow X_2$ and $g_2 : X_2 \rightarrow X_1$ are continuous operators such that $g_2 \circ g_1$ is homotopic to the identity on X_1 , then $\text{cat } X_1 \leq \text{cat } X_2$.

Definition 5. Let X be any topological space and let $H_k(X)$ denotes its k -th homology group with coefficients in \mathbb{Q} . The Poincaré polynomial $P_t(X)$ of X is defined as the following power series in t

$$P_t(X) := \sum_{k \geq 0} (\dim H_k(X)) t^k$$

Actually, if X is a compact space, we have that $\dim H_k(X) < \infty$ and this series is finite; in this case, $P_t(X)$ is a polynomial and not a formal series.

Remark 6. Let X and Y be topological spaces. If $f : X \rightarrow Y$ and $g : Y \rightarrow X$ are continuous operators such that $g \circ f$ is homotopic to the identity on X , then $P_t(Y) = P_t(X) + Z(t)$ where $Z(t)$ is a polynomial with non-negative coefficients.

These topological tools are classical and can be found, e.g., in [16] and in [5].

3. PRELIMINARY RESULTS

Using an idea in a paper of Benci and Fortunato [9] we define the map $\psi : H_0^1(\Omega) \rightarrow H_0^1(\Omega)$ defined by the equation

$$(2) \quad -\Delta\psi(u) = qu^2 \text{ in } \Omega$$

Lemma 7. *The map $\psi : H_0^1(\Omega) \rightarrow H_0^1(\Omega)$ is of class C^2 with derivatives*

$$(3) \quad \psi'(u)[\varphi] = i^*(2qu\varphi)$$

$$(4) \quad \psi''(u)[\varphi_1, \varphi_2] = i^*(2q\varphi_1\varphi_2)$$

where the operator $i_\varepsilon^* : L^{p'}, |\cdot|_{\varepsilon, p'} \rightarrow H_\varepsilon$ is the adjoint operator of the immersion operator $i_\varepsilon : H_\varepsilon \rightarrow L^p, |\cdot|_{\varepsilon, p}$.

Proof. The proof is standard. \square

Lemma 8. *The map $T : H_0^1(\Omega) \rightarrow \mathbb{R}$ given by*

$$T(u) = \int_{\Omega} u^2 \psi(u) dx$$

is a C^2 map and its first derivative is

$$T'(u)[\varphi] = 4 \int_{\Omega} \varphi u \psi(u) dx.$$

Proof. The regularity is standard. The first derivative is

$$T'(u)[\varphi] = 2 \int_{\Omega} u \varphi \psi(u) + \int_{\Omega} u^2 \psi'(u)[\varphi].$$

By (3) and (2) we have

$$\begin{aligned} 2q \int_{\Omega} u \varphi \psi(u) &= - \int_{\Omega} \Delta(\psi'(u)[\varphi]) \psi(u) = - \int_{\Omega} \psi'(u)[\varphi] \Delta\psi(u) = \\ &= \int_{\Omega} \psi'(u)[\varphi] qu^2 \end{aligned}$$

and the claim follows. \square

At this point we consider the following functional $I_\varepsilon \in C^2(H_0^1(\Omega), \mathbb{R})$.

$$(5) \quad I_\varepsilon(u) = \frac{1}{2} \|u\|_\varepsilon^2 + \frac{\omega}{4} G_\varepsilon(u) - \frac{1}{p} |u^+|_{\varepsilon, p}^p$$

where

$$G_\varepsilon(u) = \frac{1}{\varepsilon^3} \int_{\Omega} u^2 \psi(u) dx = \frac{1}{\varepsilon^3} T(u).$$

By Lemma 8 we have

$$I'_\varepsilon(u)[\varphi] = \frac{1}{\varepsilon^3} \int_{\Omega} \varepsilon^2 \nabla u \nabla \varphi + u \varphi + \omega u \psi(u) \varphi - (u^+)^{p-1} \varphi$$

$$I'_\varepsilon(u)[u] = \|u\|_\varepsilon^2 + \omega G_\varepsilon(u) - |u^+|_{\varepsilon, p}^p$$

then if u is a critical points of the functional I_ε the pair of positive functions $(u, \psi(u))$ is a solution of (1).

4. NEHARI MANIFOLD

We define the following Nehari set

$$\mathcal{N}_\varepsilon = \{u \in H_0^1(\Omega) \setminus 0 : N_\varepsilon(u) := I'_\varepsilon(u)[u] = 0\}$$

In this section we give an explicit proof of the main properties of the Nehari manifold, although standard, for the sake of completeness

Lemma 9. \mathcal{N}_ε is a C^2 manifold and $\inf_{\mathcal{N}_\varepsilon} \|u\|_\varepsilon > 0$.

Proof. If $u \in \mathcal{N}_\varepsilon$, using that $N_\varepsilon(u) = 0$, and $p > 4$ we have

$$N'_\varepsilon(u)[u] = 2\|u\|_\varepsilon^2 + 4\omega G_\varepsilon(u) - p|u^+|_{\varepsilon,p} = (2-p)\|u\|_\varepsilon + (4-p)\omega G_\varepsilon(u) < 0$$

so \mathcal{N}_ε is a C^2 manifold.

We prove the second claim by contradiction. Take a sequence $\{u_n\}_n \in \mathcal{N}_\varepsilon$ with $\|u_n\|_\varepsilon \rightarrow 0$ while $n \rightarrow +\infty$. Thus, using that $N_\varepsilon(u) = 0$,

$$\|u_n\|_\varepsilon^2 + \omega G_\varepsilon(u_n) = |u_n^+|_{p,\varepsilon}^p \leq C\|u_n\|_\varepsilon^p,$$

so

$$1 < 1 + \frac{\omega G_\varepsilon(u)}{\|u_n\|_\varepsilon} \leq C\|u_n\|_\varepsilon^{p-2} \rightarrow 0$$

and this is a contradiction. \square

Remark 10. If $u \in \mathcal{N}_\varepsilon$, then

$$\begin{aligned} I_\varepsilon(u) &= \left(\frac{1}{2} - \frac{1}{p}\right) \|u\|_\varepsilon^2 + \omega \left(\frac{1}{4} - \frac{1}{p}\right) G_\varepsilon(u) \\ &= \left(\frac{1}{2} - \frac{1}{p}\right) |u^+|_{p,\varepsilon}^p - \frac{\omega}{4} G_\varepsilon(u) \end{aligned}$$

Lemma 11. It holds Palais-Smale condition for the functional I_ε on \mathcal{N}_ε .

Proof. We start proving PS condition for I_ε . Let $\{u_n\}_n \in H_0^1(\Omega)$ such that

$$I_\varepsilon(u_n) \rightarrow c \quad |I'_\varepsilon(u_n)[\varphi]| \leq \sigma_n \|\varphi\|_\varepsilon \text{ where } \sigma_n \rightarrow 0$$

We prove that $\|u_n\|_\varepsilon$ is bounded. Suppose $\|u_n\|_\varepsilon \rightarrow \infty$. Then, by PS hypothesis

$$\frac{pI_\varepsilon(u_n) - I'_\varepsilon(u_n)[u_n]}{\|u_n\|_\varepsilon} = \left(\frac{p}{2} - 1\right) \|u_n\|_\varepsilon + \left(\frac{p}{4} - 1\right) \frac{G_\varepsilon(u_n)}{\|u_n\|_\varepsilon} \rightarrow 0$$

and this is a contradiction because $p > 4$.

At this point, up to subsequence $u_n \rightarrow u$ weakly in $H_0^1(\Omega)$ and strongly in $L^t(\Omega)$ for each $2 \leq t < 6$. Since u_n is a PS sequence

$$u_n + \omega i_\varepsilon^*(\psi(u_n)u_n) - i_\varepsilon^*((u_n^+)^{p-1}) \rightarrow 0 \text{ in } H_0^1(\Omega)$$

we have only to prove that $i_\varepsilon^*(\psi(u_n)u_n) \rightarrow i_\varepsilon^*(\psi(u)u)$ in $H_0^1(\Omega)$, then we have to prove that

$$\psi(u_n)u_n \rightarrow \psi(u)u \text{ in } L^{t'}$$

We have $|\psi(u_n)u_n - \psi(u)u|_{\varepsilon,t'} \leq |\psi(u)(u_n - u)|_{\varepsilon,t'} + |(\psi(u_n) - \psi(u))u_n|_{\varepsilon,t'}$. We get

$$\int_\Omega |\psi(u_n) - \psi(u)|^{\frac{t}{t-1}} |u_n|^{\frac{t}{t-1}} \leq \left(\int_\Omega |\psi(u_n) - \psi(u)|^t \right)^{\frac{1}{t-1}} \left(\int_\Omega |u_n|^{\frac{t}{t-2}} \right)^{\frac{t-2}{t-1}} \rightarrow 0,$$

thus we can conclude easily.

Now we prove PS condition for the constrained functional. Let $\{u_n\}_n \in \mathcal{N}_\varepsilon$ such that

$$I_\varepsilon(u_n) \rightarrow c \\ |I'_\varepsilon(u_n)[\varphi] - \lambda_n N'(u_n)[\varphi]| \leq \sigma_n \|\varphi\|_\varepsilon \quad \text{with } \sigma_n \rightarrow 0$$

In particular $I'_\varepsilon(u_n) \left[\frac{u_n}{\|u_n\|_\varepsilon} \right] - \lambda_n N'(u_n) \left[\frac{u_n}{\|u_n\|_\varepsilon} \right] \rightarrow 0$. Then

$$\lambda_n \left\{ (p-2) \|u_n\|_\varepsilon + (p-4) \omega \frac{G_\varepsilon(u_n)}{\|u_n\|_\varepsilon} \right\} \rightarrow 0$$

thus $\lambda_n \rightarrow 0$ because $p > 4$. Since $N'(u_n) = u_n - i_\varepsilon^*(4\omega\psi(u_n)u_n) - pi_\varepsilon^*(|u_n^+|^{p-1})$ is bounded we obtain that $\{u_n\}_n$ is a PS sequence for the free functional I_ε , and we get the claim \square

Lemma 12. *For all $w \in H_0^1(\Omega)$ such that $|w^+|_{\varepsilon,p} = 1$ there exists a unique positive number $t_\varepsilon = t_\varepsilon(w)$ such that $t_\varepsilon(w)w \in \mathcal{N}_\varepsilon$.*

Proof. We define, for $t > 0$

$$H(t) = I_\varepsilon(tw) = \frac{1}{2}t^2\|w\|_\varepsilon^2 + \frac{t^4}{4}\omega G_\varepsilon(w) - \frac{t^p}{p}.$$

Thus

$$(6) \quad H'(t) = t(\|w\|_\varepsilon^2 + t^2\omega G_\varepsilon(w) - t^{p-2})$$

$$(7) \quad H''(t) = \|w\|_\varepsilon^2 + 3t^2\omega G_\varepsilon(w) - (p-1)t^{p-2}$$

By (6) there exists $t_\varepsilon > 0$ such that $H'(t_\varepsilon) = 0$. Moreover, by (6), (7) and because $p > 4$ we have that $H''(t_\varepsilon) < 0$, so t_ε is unique. \square

5. MAIN INGREDIENT OF THE PROOF

We sketch the proof of Theorem 1. First of all, since the functional $I_\varepsilon \in C^2$ is bounded below and satisfies PS condition on the complete C^2 manifold \mathcal{N}_ε , we have, by well known results, that I_ε has at least $\text{cat } I_\varepsilon^d$ critical points in the sublevel

$$I_\varepsilon^d = \{u \in H^1 : I_\varepsilon(u) \leq d\}.$$

We prove that, for ε and δ small enough, it holds

$$\text{cat } \Omega \leq \text{cat } (\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta})$$

where

$$m_\infty := \inf_{\mathcal{N}_\infty} \frac{1}{2} \int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx - \frac{1}{p} \int_{\mathbb{R}^3} |v|^p dx \\ \mathcal{N}_\infty = \left\{ v \in H^1(\mathbb{R}^3) \setminus \{0\} : \int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx = \int_{\mathbb{R}^3} |v|^p dx \right\}.$$

To get the inequality $\text{cat } \Omega \leq \text{cat } (\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta})$ we build two continuous operators

$$\begin{aligned} \Phi_\varepsilon &: \Omega^- \rightarrow \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta} \\ \beta &: \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta} \rightarrow \Omega^+. \end{aligned}$$

where

$$\begin{aligned} \Omega^- &= \{x \in \Omega : d(x, \partial\Omega) < r\} \\ \Omega^+ &= \{x \in \mathbb{R}^3 : d(x, \partial\Omega) < r\} \end{aligned}$$

with r small enough so that $\text{cat}(\Omega^-) = \text{cat}(\Omega^+) = \text{cat}(\Omega)$.

Following an idea in [7], we build these operators Φ_ε and β such that $\beta \circ \Phi_\varepsilon : \Omega^- \rightarrow \Omega^+$ is homotopic to the immersion $i : \Omega^- \rightarrow \Omega^+$. By the properties of Lusternik Schinerlmann category we have

$$\text{cat } \Omega \leq \text{cat } (\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta})$$

which ends the proof of Theorem 1.

Concerning Theorem 2, we can re-state classical results contained in [5, 8] in the following form.

Theorem 13. *Let I_ε be the functional (5) on $H^1(\Omega)$ and let K_ε be the set of its critical points. If all its critical points are non-degenerate then*

$$(8) \quad \sum_{u \in K_\varepsilon} t^{\mu(u)} = tP_t(\Omega) + t^2(P_t(\Omega) - 1) + t(1+t)Q(t)$$

where $Q(t)$ is a polynomial with non-negative integer coefficients and $\mu(u)$ is the Morse index of the critical point u .

By Remark 6 and by means of the maps Φ_ε and β we have that

$$(9) \quad P_t(\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta}) = P_t(\Omega) + Z(t)$$

where $Z(t)$ is a polynomial with non-negative coefficients. Provided that $\inf_\varepsilon m_\varepsilon =: \alpha > 0$, because $\lim_{\varepsilon \rightarrow 0} m_\varepsilon = m_\infty$ (see 20), we have the following relations [5, 8]

$$(10) \quad P_t(I_\varepsilon^{m_\infty + \delta}, I_\varepsilon^{\alpha/2}) = tP_t(\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta})$$

$$(11) \quad P_t(H_0^1(\Omega), I_\varepsilon^{m_\infty + \delta}) = t(P_t(I_\varepsilon^{m_\infty + \delta}, I_\varepsilon^{\alpha/2}) - t)$$

$$(12) \quad \sum_{u \in K_\varepsilon} t^{\mu(u)} = P_t(H_0^1(\Omega), I_\varepsilon^{m_\infty + \delta}) + P_t(I_\varepsilon^{m_\infty + \delta}, I_\varepsilon^{\alpha/2}) + (1+t)\tilde{Q}(t)$$

where $\tilde{Q}(t)$ is a polynomial with non-negative integer coefficients. Hence, by (9), (10), (11), (12) we obtain (8). At this point, evaluating equation (8) for $t = 1$ we obtain the claim of Theorem 2

6. THE MAP Φ_ε

For every $\xi \in \Omega^-$ we define the function

$$W_{\xi, \varepsilon}(x) = U_\varepsilon(x - \xi)\chi(|x - \xi|)$$

where $\chi : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ where $\chi \equiv 1$ for $t \in [0, r/2)$, $\chi \equiv 0$ for $t > r$ and $|\chi'(t)| \leq 2/r$.

We can define a map

$$\begin{aligned} \Phi_\varepsilon & : \Omega^- \rightarrow \mathcal{N}_\varepsilon \\ \Phi_\varepsilon(\xi) & = t_\varepsilon(W_{\xi, \varepsilon})W_{\xi, \varepsilon} \end{aligned}$$

Remark 14. We have that the following limits hold uniformly with respect to $\xi \in \Omega$

$$\begin{aligned} \|W_{\varepsilon, \xi}\|_\varepsilon & \rightarrow \|U\|_{H^1(\mathbb{R}^3)} \\ |W_{\varepsilon, \xi}|_{\varepsilon, t} & \rightarrow \|U\|_{L^t(\mathbb{R}^3)} \text{ for all } 2 \leq t \leq 6 \end{aligned}$$

Lemma 15. *There exists $\bar{\varepsilon} > 0$ and a constant $c > 0$ such that*

$$G_\varepsilon(W_{\varepsilon, \xi}) = \frac{1}{\varepsilon^3} \int_\Omega qW_{\varepsilon, \xi}^2(x)\psi(W_{\varepsilon, \xi})dx < c\varepsilon^2$$

Proof. It holds

$$\begin{aligned} \|\psi(W_{\varepsilon,\xi})\|_{H_0^1(\Omega)}^2 &= \int_{\Omega} q W_{\varepsilon,\xi}^2(x) \psi(W_{\varepsilon,\xi}) dx \leq q \|\psi(W_{\varepsilon,\xi})\|_{L^6(\Omega)} \left(\int_{\Omega} W_{\varepsilon,\xi}^{12/5} dx \right)^{5/6} \\ &\leq c \|\psi(W_{\varepsilon,\xi})\|_{H_0^1(\Omega)} \left(\frac{1}{\varepsilon^3} \int_{\Omega} W_{\varepsilon,\xi}^{12/5} dx \right)^{5/6} \varepsilon^{5/2} \end{aligned}$$

By Remark 14 we have that $\|\psi(W_{\varepsilon,\xi})\|_{H_0^1(\Omega)} \leq \varepsilon^{5/2}$ and the claim follows by applying again Cauchy Schwartz inequality. \square

Proposition 16. *For all $\varepsilon > 0$ the map Φ_{ε} is continuous. Moreover for any $\delta > 0$ there exists $\varepsilon_0 = \varepsilon_0(\delta)$ such that, if $\varepsilon < \varepsilon_0$ then $I_{\varepsilon}(\Phi_{\varepsilon}(\xi)) < m_{\infty} + \delta$.*

Proof. It is easy to see that Φ_{ε} is continuous because $t_{\varepsilon}(w)$ depends continuously on $w \in H_0^1$.

At this point we prove that $t_{\varepsilon}(W_{\varepsilon,\xi}) \rightarrow 1$ uniformly with respect to $\xi \in \Omega$. In fact, by Lemma 12 $t_{\varepsilon}(W_{\varepsilon,\xi})$ is the unique solution of

$$\|W_{\varepsilon,\xi}\|_{\varepsilon}^2 + t^2 \omega G_{\varepsilon}(W_{\varepsilon,\xi}) - t^{p-2} |W_{\varepsilon,\xi}|_{\varepsilon,p}^p = 0.$$

By Remark 14 and Lemma 15 we have the claim.

Now, we have

$$I_{\varepsilon}(t_{\varepsilon}(W_{\varepsilon,\xi})W_{\varepsilon,\xi}) = \left(\frac{1}{2} - \frac{1}{p} \right) \|W_{\varepsilon,\xi}\|_{\varepsilon}^2 t_{\varepsilon}^2 + \omega \left(\frac{1}{4} - \frac{1}{p} \right) t_{\varepsilon}^4 G_{\varepsilon}(W_{\varepsilon,\xi})$$

Again, by Remark 14 and Lemma 15 we have

$$I_{\varepsilon}(t_{\varepsilon}(W_{\varepsilon,\xi})W_{\varepsilon,\xi}) \rightarrow \left(\frac{1}{2} - \frac{1}{p} \right) \|U\|_{H^1(\mathbb{R}^3)}^2 = m_{\infty}$$

that concludes the proof. \square

Remark 17. We set

$$m_{\varepsilon} = \inf_{\mathcal{N}_{\varepsilon}} I_{\varepsilon}.$$

By Proposition 16 we have that

$$(13) \quad \limsup_{\varepsilon \rightarrow 0} m_{\varepsilon} \leq m_{\infty}.$$

7. THE MAP β

For any $u \in \mathcal{N}_{\varepsilon}$ we can define a point $\beta(u) \in \mathbb{R}^3$ by

$$\beta(u) = \frac{\int_{\Omega} x |u^+|^p dx}{\int_{\Omega} |u^+|^p dx}.$$

The function β is well defined in $\mathcal{N}_{\varepsilon}$ because, if $u \in \mathcal{N}_{\varepsilon}$, then $u^+ \neq 0$.

We have to prove that, if $u \in \mathcal{N}_{\varepsilon} \cap I_{\varepsilon}^{m_{\infty} + \delta}$ then $\beta(u) \in \Omega^+$.

Let us consider partitions of Ω . For a given $\varepsilon > 0$ we say that a finite partition $\mathcal{P}_{\varepsilon} = \{P_j^{\varepsilon}\}_{j \in \Lambda_{\varepsilon}}$ of Ω is a “good” partition if: for any $j \in \Lambda_{\varepsilon}$ the set P_j^{ε} is closed; $P_i^{\varepsilon} \cap P_j^{\varepsilon} \subset \partial P_i^{\varepsilon} \cap \partial P_j^{\varepsilon}$ for any $i \neq j$; there exist $r_1(\varepsilon), r_2(\varepsilon) > 0$ such that there are points $q_j^{\varepsilon} \in P_j^{\varepsilon}$ for which $B(q_j^{\varepsilon}, \varepsilon) \subset P_j^{\varepsilon} \subset B(q_j^{\varepsilon}, r_2(\varepsilon)) \subset B(q_j^{\varepsilon}, r_1(\varepsilon))$, with $r_1(\varepsilon) \geq r_2(\varepsilon) \geq C\varepsilon$ for some positive constant C ; lastly, there exists a finite number

$\nu \in \mathbb{N}$ such that every $x \in \Omega$ is contained in at most ν balls $B(q_j^\varepsilon, r_1(\varepsilon))$, where ν does not depends on ε .

Lemma 18. *There exists a constant $\gamma > 0$ such that, for any $\delta > 0$ and for any $\varepsilon < \varepsilon_0(\delta)$ as in Proposition 16, given any “good” partition $\mathcal{P}_\varepsilon = \{P_j^\varepsilon\}_j$ of the domain Ω and for any function $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta}$ there exists, for an index \bar{j} a set $P_{\bar{j}}^\varepsilon$ such that*

$$\frac{1}{\varepsilon^3} \int_{P_{\bar{j}}^\varepsilon} |u^+|^p dx \geq \gamma.$$

Proof. Taking in account that $I'(u)[u] = 0$ we have

$$\begin{aligned} \|u\|_\varepsilon^2 &= |u^+|_{\varepsilon,p}^p - \frac{1}{\varepsilon^3} \int_\Omega \omega u^2 \psi(u) \leq |u^+|_{\varepsilon,p}^p = \sum_j \frac{1}{\varepsilon^3} \int_{P_j} |u^+|^p \\ &= \sum_j |u_j^+|_{\varepsilon,p}^p = \sum_j |u_j^+|_{\varepsilon,p}^{p-2} |u_j^+|_{\varepsilon,p}^2 \leq \max_j \{|u_j^+|_{\varepsilon,p}^{p-2}\} \sum_j |u_j^+|_{\varepsilon,p}^2 \end{aligned}$$

where u_j^+ is the restriction of the function u^+ on the set P_j .

At this point, arguing as in [6, Lemma 5.3], we prove that there exists a constant $C > 0$ such that

$$\sum_j |u_j^+|_{\varepsilon,p}^2 \leq C \nu \|u^+\|_\varepsilon^2,$$

thus

$$\max_j \{|u_j^+|_{\varepsilon,p}^{p-2}\} \geq \frac{1}{C \nu}$$

that concludes the proof. \square

Proposition 19. *For any $\eta \in (0, 1)$ there exists $\delta_0 < m_\infty$ such that for any $\delta \in (0, \delta_0)$ and any $\varepsilon \in (0, \varepsilon_0(\delta))$ as in Proposition 16, for any function $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta}$ we can find a point $q = q(u) \in \Omega$ such that*

$$\frac{1}{\varepsilon^3} \int_{B(q,r/2)} (u^+)^p > (1 - \eta) \frac{2p}{p-2} m_\infty.$$

Proof. First, we prove the proposition for $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + 2\delta}$.

By contradiction, we assume that there exists $\eta \in (0, 1)$ such that we can find two sequences of vanishing real number δ_k and ε_k and a sequence of functions $\{u_k\}_k$ such that $u_k \in \mathcal{N}_{\varepsilon_k}$,

(14)

$$m_{\varepsilon_k} \leq I_{\varepsilon_k}(u_k) = \left(\frac{1}{2} - \frac{1}{p}\right) \|u_k\|_{\varepsilon_k}^2 + \omega \left(\frac{1}{4} - \frac{1}{p}\right) G_{\varepsilon_k}(u_k) \leq m_{\varepsilon_k} + 2\delta_k \leq m_\infty + 3\delta_k$$

for k large enough (see Remark 17), and, for any $q \in \Omega$,

$$\frac{1}{\varepsilon_k^3} \int_{B(q,r/2)} (u_k^+)^p \leq (1 - \eta) \frac{2p}{p-2} m_\infty.$$

By Ekeland principle and by definition of $\mathcal{N}_{\varepsilon_k}$ we can assume

$$(15) \quad |I'_{\varepsilon_k}(u_k)[\varphi]| \leq \sigma_k \|\varphi\|_{\varepsilon_k} \text{ where } \sigma_k \rightarrow 0.$$

By Lemma 18 there exists a set $P_k^{\varepsilon_k} \in \mathcal{P}_{\varepsilon_k}$ such that

$$\frac{1}{\varepsilon_k^3} \int_{P_k^{\varepsilon_k}} |u_k^+|^p dx \geq \gamma.$$

We choose a point $q_k \in P_k^{\varepsilon_k}$ and we define, for $z \in \Omega_{\varepsilon_k} := \frac{1}{\varepsilon_k}(\Omega - q_k)$

$$w_k(z) = u_k(\varepsilon_k z + q_k) = u_k(x).$$

We have that $w_k \in H_0^1(\Omega_{\varepsilon_k}) \subset H^1(\mathbb{R}^3)$. By equation (14) we have

$$\|w_k\|_{H^1(\mathbb{R}^3)}^2 = \|u_k\|_{\varepsilon_k}^2 \leq C.$$

So $w_k \rightarrow w$ weakly in $H^1(\mathbb{R}^3)$ and strongly in $L_{\text{loc}}^t(\mathbb{R}^3)$.

We set $\psi(u_k)(x) := \psi_k(x) = \psi_k(\varepsilon_k z + q_k) := \tilde{\psi}_k(z)$ where $x \in \Omega$ and $z \in \Omega_{\varepsilon_k}$. It is easy to verify that

$$-\Delta_z \tilde{\psi}_k(z) = \varepsilon_k^2 q w_k^2(z).$$

With abuse of language we set

$$\tilde{\psi}_k(z) = \psi(\varepsilon_k w_k).$$

Thus

$$\begin{aligned} I_{\varepsilon_k}(u_k) &= \frac{1}{2} \|u_k\|_{\varepsilon_k}^2 - \frac{1}{p} |u_k^+|_{\varepsilon_k, p}^p + \frac{\omega}{4} \frac{1}{\varepsilon_k^3} \int_{\Omega} q u_k^2 \psi(u_k) = \\ (16) \quad &= \frac{1}{2} \|w_k\|_{H^1(\mathbb{R}^3)}^2 - \frac{1}{p} \|w_k^+\|_{L^p(\mathbb{R}^3)}^p + \frac{\omega}{4} \int_{\Omega_{\varepsilon_k}} q w_k^2 \psi(\varepsilon_k w_k) = \\ &= \frac{1}{2} \|w_k\|_{H^1(\mathbb{R}^3)}^2 - \frac{1}{p} \|w_k^+\|_{L^p(\mathbb{R}^3)}^p + \varepsilon_k^2 \frac{\omega}{4} \int_{\mathbb{R}^3} q w_k^2 \psi(w_k) := E_{\varepsilon_k}(w_k) \end{aligned}$$

By definition of $E_{\varepsilon_k} : H^1(\mathbb{R}^3) \rightarrow \mathbb{R}$, we get $E_{\varepsilon_k}(w_k) \rightarrow m_{\infty}$.

Given any $\varphi \in C_0^{\infty}(\mathbb{R}^3)$ we set $\varphi(x) = \varphi(\varepsilon_k z + q_k) := \tilde{\varphi}_k(z)$. For k large enough we have that $\text{supp } \tilde{\varphi}_k \subset \Omega$ and, by (15), that $E'_{\varepsilon_k}(w_k)[\varphi] = I'_{\varepsilon_k}(u_k)[\tilde{\varphi}_k] \rightarrow 0$. Moreover, by definition of E_{ε_k} and by Lemma 8 we have

$$\begin{aligned} E'_{\varepsilon_k}(w_k)[\varphi] &= \langle w_k, \varphi \rangle_{H^1(\mathbb{R}^3)} - \int_{\mathbb{R}^3} |w_k^+|^{p-1} \varphi + \omega \varepsilon_k^2 \int_{\mathbb{R}^3} q w_k \psi(w_k) \varphi + \\ &\rightarrow \langle w, \varphi \rangle_{H^1(\mathbb{R}^3)} - \int_{\mathbb{R}^3} |w^+|^{p-1} \varphi. \end{aligned}$$

Thus w is a weak solution of

$$-\Delta w + w = (w^+)^{p-1} \text{ on } \mathbb{R}^3.$$

By Lemma 18 and by the choice of q_k we have that $w \neq 0$, so $w > 0$.

Arguing as in (16), and using that $u_k \in \mathcal{N}_{\varepsilon_k}$ we have

$$\begin{aligned} (17) \quad I_{\varepsilon_k}(u_k) &= \left(\frac{1}{2} - \frac{1}{p} \right) \|u_k\|_{\varepsilon_k}^2 + \omega \left(\frac{1}{4} - \frac{1}{p} \right) \frac{1}{\varepsilon_k^3} \int_{\Omega} q u_k^2 \psi(u_k) \\ &= \left(\frac{1}{2} - \frac{1}{p} \right) \|w_k\|_{H^1(\mathbb{R}^3)}^2 + \varepsilon_k^2 \omega \left(\frac{1}{4} - \frac{1}{p} \right) \int_{\mathbb{R}^3} q w_k^2 \psi(w_k) \rightarrow m_{\infty} \end{aligned}$$

and

$$\begin{aligned} (18) \quad I_{\varepsilon_k}(u_k) &= \left(\frac{1}{2} - \frac{1}{p} \right) |u_k^+|_{p, \varepsilon_k}^p - \frac{\omega}{4} \frac{1}{\varepsilon_k^3} \int_{\Omega} q u_k^2 \psi(u_k) \\ &= \left(\frac{1}{2} - \frac{1}{p} \right) |w_k^+|_p^p - \varepsilon_k^2 \frac{\omega}{4} \int_{\mathbb{R}^3} q w_k^2 \psi(w_k) \rightarrow m_{\infty}. \end{aligned}$$

So, by (17) we have that $\|w\|_{H^1(\mathbb{R}^3)}^2 = \frac{2p}{p-2} m_{\infty}$ and that $\left(\frac{1}{2} - \frac{1}{p} \right) \|w_k\|_{H^1(\mathbb{R}^3)}^2 \rightarrow m_{\infty}$ and we conclude that $w_k \rightarrow w$ strongly in $H^1(\mathbb{R}^3)$.

Given $T > 0$, by the definition of w_k we get, for k large enough

$$\begin{aligned} |w_k^+|_{L^p(B(0,T))}^p &= \frac{1}{\varepsilon_k^3} \int_{B(q_k, \varepsilon_k T)} |u_k^+|^p dx \leq \frac{1}{\varepsilon_k^3} \int_{B(q_k, r/2)} |u_k^+|^p dx \\ (19) \quad &\leq (1 - \eta) \frac{2p}{p-2} m_\infty. \end{aligned}$$

Then we have the contradiction. In fact, by (18) we have $\left(\frac{1}{2} - \frac{1}{p}\right) |w_k^+|_p^p \rightarrow m_\infty$ and this contradicts (19). At this point we have proved the claim for $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\varepsilon+2\delta}$. Now, by the thesis for $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\varepsilon+2\delta}$ and by (18) we have

$$I_{\varepsilon_k}(u_k) = \left(\frac{1}{2} - \frac{1}{p}\right) |u_k^+|_{p, \varepsilon_k}^p + O(\varepsilon^2) \geq (1 - \eta) m_\infty + O(\varepsilon^2)$$

and, passing to the limit,

$$\liminf_{k \rightarrow \infty} m_{\varepsilon_k} \geq m_\infty.$$

This, combined by (13) gives us that

$$(20) \quad \lim_{\varepsilon \rightarrow 0} m_\varepsilon = m_\infty.$$

Hence, when ε, δ are small enough, $\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta} \subset \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\varepsilon+2\delta}$ and the general claim follows. \square

Proposition 20. *There exists $\delta_0 \in (0, m_\infty)$ such that for any $\delta \in (0, \delta_0)$ and any $\varepsilon \in (0, \varepsilon(\delta_0))$ (see Proposition 16), for every function $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$ it holds $\beta(u) \in \Omega^+$. Moreover the composition*

$$\beta \circ \Phi_\varepsilon : \Omega^- \rightarrow \Omega^+$$

is homotopic to the immersion $i : \Omega^- \rightarrow \Omega^+$

Proof. By Proposition 19, for any function $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$, for any $\eta \in (0, 1)$ and for ε, δ small enough, we can find a point $q = q(u) \in \Omega$ such that

$$\frac{1}{\varepsilon^3} \int_{B(q, r/2)} (u^+)^p > (1 - \eta) \frac{2p}{p-2} m_\infty.$$

Moreover, since $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$ we have

$$I_\varepsilon(u) = \left(\frac{p-2}{2p}\right) |u^+|_{p, \varepsilon}^p - \frac{\omega}{4} \frac{1}{\varepsilon^3} \int_\Omega q u^2 \psi(u) \leq m_\infty + \delta.$$

Now, arguing as in Lemma 15 we have that

$$\|\psi(u)\|_{H^1(\Omega)}^2 = q \int_\Omega \psi(u) u^2 \leq C \|\psi(u)\|_{H^1(\Omega)} \left(\int_\Omega u^{12/5} \right)^{5/6},$$

so $\|\psi(u)\|_{H^1(\Omega)} \leq \left(\int_\Omega u^{12/5} \right)^{5/6}$, then

$$\begin{aligned} \frac{1}{\varepsilon^3} \int \psi(u) u^2 &\leq \frac{1}{\varepsilon^3} \|\psi\|_{H^1(\Omega)} \left(\int_\Omega u^{12/5} \right)^{5/6} \leq C \frac{1}{\varepsilon^3} \left(\int_\Omega u^{12/5} \right)^{5/3} \\ &\leq C \varepsilon^2 |u|_{12/5, \varepsilon}^4 \leq C \varepsilon^2 \|u\|_\varepsilon^4 \leq C \varepsilon^2 \end{aligned}$$

because $\|u\|_\varepsilon$ is bounded since $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$.

Hence, provided we choose $\varepsilon(\delta_0)$ small enough, we have

$$\left(\frac{p-2}{2p}\right) |u^+|_{p,\varepsilon}^p \leq m_\infty + 2\delta_0.$$

So,

$$\frac{\frac{1}{\varepsilon^3} \int_{B(q,r/2)} (u^+)^p}{|u^+|_{p,\varepsilon}^p} > \frac{1-\eta}{1+2\delta_0/m_\infty}$$

Finally,

$$\begin{aligned} |\beta(u) - q| &\leq \frac{\left| \frac{1}{\varepsilon^3} \int_{\Omega} (x-q)(u^+)^p \right|}{|u^+|_{p,\varepsilon}^p} \\ &\leq \frac{\left| \frac{1}{\varepsilon^3} \int_{B(q,r/2)} (x-q)(u^+)^p \right|}{|u^+|_{p,\varepsilon}^p} + \frac{\left| \frac{1}{\varepsilon^3} \int_{\Omega \setminus B(q,r/2)} (x-q)(u^+)^p \right|}{|u^+|_{p,\varepsilon}^p} \\ &\leq \frac{r}{2} + 2\text{diam}(\Omega) \left(1 - \frac{1-\eta}{1+2\delta_0/m_\infty} \right), \end{aligned}$$

so, choosing η , δ_0 and $\varepsilon(\delta_0)$ small enough we proved the first claim. The second claim is standard. \square

REFERENCES

- [1] A. Ambrosetti, D. Ruiz, *Multiple bound states for the Schroedinger-Poisson problem*, Commun. Contemp. Math. **10** (2008) 391–404
- [2] A. Azzollini, P. D’Avenia, A. Pomponio, *On the Schroedinger-Maxwell equations under the effect of a general nonlinear term*, Ann. Inst. H. Poincaré Anal. Non Linéaire **27** (2010), no. 2, 779–791
- [3] A. Azzollini, A. Pomponio, *Ground state solutions for the nonlinear Schroedinger-Maxwell equations*, J. Math. Anal. Appl. **345** (2008) no. 1, 90–108
- [4] J. Bellazzini, L. Jeanjean, T. Luo, *Existence and instability of standing waves with prescribed norm for a class of Schroedinger-Poisson equations* in press on Proc. London Math. Soc. (arXiv <http://arxiv.org/abs/1111.4668>)
- [5] V. Benci, *Introduction to Morse theory: A new approach*, in: Topological Nonlinear Analysis, in: Progr. Nonlinear Differential Equations Appl., vol. 15, Birkhauser Boston, Boston, MA, 1995, pp. 37–177.
- [6] V. Benci, C. Bonanno, A.M. Micheletti, *On the multiplicity of solutions of a nonlinear elliptic problem on Riemannian manifolds*, Journal of Functional Analysis **252** (2007) 464–489
- [7] V. Benci, G. Cerami, *The effect of the domain topology on the number of positive solutions of nonlinear elliptic problems*, Arch. Ration. Mech. Anal. **114** (1991) 79–93.
- [8] V. Benci, G. Cerami, *Multiple positive solutions of some elliptic problems via the Morse theory and the domain topology*, Calc. Var. Partial Differential Equations **2** (1994) 29–48.
- [9] V. Benci, D. Fortunato, *An eigenvalue problem for the Schroedinger-Maxwell equations*, Topol. Methods Nonlinear Anal. **11** (1998), no. 2, 283–293
- [10] T. D’Aprile, D. Mugnai, *Solitary waves for nonlinear Klein-Gordon-Maxwell and Schroedinger-Maxwell equations*, Proc. Roy. Soc. Edinburgh Sect. A **134** (2004), no. 5, 893–906.
- [11] T. D’Aprile, J. Wei, *Layered solutions for a semilinear elliptic system in a ball*, J. Differential Equations **226** (2006) , no. 1, 269–294.
- [12] T. D’Aprile, J. Wei, *Clustered solutions around harmonic centers to a coupled elliptic system*, Ann. Inst. H. Poincaré Anal. Non Linéaire **24** (2007), no. 4, 605–628.
- [13] M. Ghimenti, A.M. Micheletti, *Number and profile of low energy solutions for singularly perturbed Klein Gordon Maxwell systems on a Riemannian manifold*, work in preparation.
- [14] I. Ianni, G. Vaira, *On concentration of positive bound states for the Schroedinger-Poisson problem with potentials*, Adv. Nonlinear Stud. **8** (2008), no. 3, 573–595.

- [15] Kikuchi, *On the existence of solutions for a elliptic system related to the Maxwell-Schroedinger equations*, Non- linear Anal. **67** (2007) 1445–1456.
- [16] R. S. Palais, *Homotopy theory of infinite dimensional manifolds*, Topology **5** (1966), 1–16.
- [17] L. Pisani, G. Siciliano, *Note on a Schroedinger-Poisson system in a bounded domain*, Appl. Math. Lett. **21** (2008), no. 5, 521–528.
- [18] D. Ruiz, *Semiclassical states for coupled Schroedinger-Maxwell equations: Concentration around a sphere*, Math. Models Methods Appl. Sci. **15** (2005), no. 1, 141–164.
- [19] G. Siciliano, *Multiple positive solutions for a Schroedinger-Poisson-Slater system*, J. Math. Anal. Appl. **365** (2010), no. 1, 288–299.
- [20] Z. Wang, H.S. Zhou, *Positive solution for a nonlinear stationary Schroedinger-Poisson system in \mathbb{R}^3* , Discrete Contin. Dyn. Syst. **18** (2007) 809–816.

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